THE LYAPUNOV SPCTRUM AS THE NEWTON METHOD.

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ABSTRACT. For a class of dynamical systems, the cookie-cutter maps, we prove that the Lyapunov spectrum coincides with the map given by the Newton-Raphson method applied to the derivative of the pressure function.

1. Introduction

In this note we establish a relation between the *Newton-Raphson* method for finding roots of an equation and an important function in the dimension theory of dynamical systems, namely the *Lyapunov spectrum*.

The dimension theory of dynamical systems has received a great deal of attention over the last years (see [Ba, Pe]). Multifractal analysis is a sub-area of dimension theory devoted to study the complexity of level sets determined by invariant local quantities. Usually, the geometry of the level sets is complicated and in order to quantify its size the Hausdorff dimension is used. In this note we will focus our attention in one of these local quantities, namely the Lyapunov exponent. Let $T:I\to I$ be a piecewise differentiable map defined on the unit interval I=[0,1]. The Lyapunov exponent of T at the point $x\in I$, is defined by

$$\lambda_T(x) = \lambda(x) = \lim_{n \to \infty} \frac{1}{n} \log |(T^n)'(x)|,$$

whenever the limit exists. It is a dynamical quantity that measures the exponential rate of divergence of infinitesimally close orbits. It is possible for the Lyapunov exponent to attain a wide range of values. In order to understand the complexity of the induced level sets we study the Lyapunov spectrum, which is the function defined by

$$L(\alpha) := \dim_H(\{x \in I : \lambda(x) = \alpha\}).$$

Here, \dim_H denotes the Hausdorff dimension of a set (see [Fa, Chapter 2] for a precise definition). This function can be described using thermodynamical formalism, in particular it can be written as a function of the topological pressure $P(-t \log |T'|)$, see Section 3 for precise definitions and statements.

We will relate the Lyapunov spectrum with the function used to apply the Newton-Raphson method. Given a real function $f: \mathbb{R} \to \mathbb{R}$, under certain assumptions, the iterates of the Newton function $N_f(x)$ converge to a root of the equation f(x) = 0 (for precise statements see section 2.2). The study of this functions and its dynamic behaviour has attracted a great deal of attention (see [Mi]).

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The main result in this note is that, for a certain class of dynamical systems, the Lyapunov spectrum equals the Newton function of the topological pressure $P(-t \log |T'|)$

$$L(-P'(-t \log |T'|)) = N_P(t).$$

2. The Legendre transform and the Newton method

Let $f: \mathbb{R} \to \mathbb{R}$ be a strictly decreasing convex function of class C^2 . In what follows we will define its corresponding Legendre transform and the associated Newton map. Moreover, we will show how are they related.

2.1. **Legendre transform.** Let $t' \in \mathbb{R}$, the line $S_{\alpha} = \alpha(t - t') + f(t')$ is called a *support line* at t' if it always lie below the graph of f. There exists an interval of the form $[\alpha_{\min}, \alpha_{\max}]$ (this interval need not to be bounded) such that if $\alpha \in [\alpha_{\min}, \alpha_{\max}]$ the graph of f has a support line S_{α} of slope equal to $-\alpha$ and for such α the support line is unique. The *Legendre transform* of f is the function $F: [\alpha_{\min}, \alpha_{\max}] \to \mathbb{R}$ defined as the value of the intercept of S_{α} with the vertical axis. That is

(1)
$$F(\alpha) = \inf\{f(t) + \alpha t : t \in \mathbb{R}\}.$$

It turns out that if $t_{\alpha} \in \mathbb{R}$ is the unique point for which $f'(t_{\alpha}) = -\alpha$ then

(2)
$$F(\alpha) = f(t_{\alpha}) + \alpha t_{\alpha}.$$

Let us remark that in this simple setting the support line S_{α} is the tangent to the graph of f at the point $(t_{\alpha}, f(t_{\alpha}))$ and that the interval $[\alpha_{\min}, \alpha_{\max}]$ is uniquely determined by the range of the derivative of f. We stress that this form of the Legendre transform is not the most commonly used in physics (see for example [A, p.61]).

Geometrically, the Legendre transform is the intercept of the tangent line of f with the vertical axis.

2.2. The Newton-Raphson method. The Newton-Raphson method is an algorithmic method that provides a sequences of approximations to the root of the equation f(t) = 0. The algorithm is based on the study of the dynamics of the so called *Newton map*. This map is defined by

$$(3) N_f(t) = t - \frac{f(t)}{f'(t)}.$$

Note that since the function f is strictly decreasing the Newton map is defined over the real numbers (the function f has no critical values). It turns out that a point $d \in \mathbb{R}$ is a root of f if and only if it is a fixed point for N_f . Moreover, since

(4)
$$N'_f(t) = \frac{f(t)f''(t)}{(f'(t))^2},$$

we have that N'(d) = 0. In particular, the point d is an attractor for N_f and iterations of the Newton map converge to the root d. That is, for any $x \in \mathbb{R}$ the sequence $\{x, N_f(x), N_f^2(x), \dots, N_f^n(x), \dots\}$ converges to the root.

Geometrically, the procedure involved in this algorithm is to consider the tangent to the graph of f at the point (t, f(t)) and define the Newton transformation as the intercept of the tangent with the horizontal axis.

2.3. A relation between the Legendre transform and the Newton method. In virtue of the above we have that in order to compute the Legendre transform and the Newton function of f we need to compute the tangent to f.

Lemma 2.1. Let $f: \mathbb{R} \to \mathbb{R}$ be a strictly decreasing convex function of class C^2 and let $t_{\alpha} \in \mathbb{R}$ be a point such that $f'(t_{\alpha}) = -\alpha$. Then

(5)
$$N_f(t_\alpha) = \frac{1}{\alpha} F(\alpha).$$

Proof. Note that since $t_{\alpha} \in \mathbb{R}$ is a point such that $f'(t_{\alpha}) = -\alpha$ then in virtue of equation (3) the Newton map is given by

$$N_f(t_\alpha) = t_\alpha + \frac{f(t_\alpha)}{\alpha}$$

and the Legendre transform (see equation (2)) is given by

$$F(\alpha) = f(t_{\alpha}) + \alpha t_{\alpha}.$$

The result now follows.

3. Thermodynamic formalism

A large class of interesting dynamical systems have many invariant measures. It is, therefore, an important problem to find criteria to choose *relevant* invariant measures. Thermodynamic formalism is a set of ideas and techniques which derive from statistical mechanics and that was brought into dynamics in the early seventies by Ruelle and Sinai among others (see the books [K, Ru2, Wa] for a detailed and nice exposition of the subject). This formalism provides procedures for the choice of interesting measures. In this note we will consider the following class of dynamical systems. Given a pairwise disjoint finite family of closed intervals I_1, \ldots, I_n contained in [0,1] we say that a map:

$$T: \bigcup_{i=1}^{n} I_i \to [0,1]$$

is a cookie-cutter map with n branches if the following holds:

- (1) $T(I_i) = [0,1]$ for every $i \in \{1,\ldots,n\}$,
- (2) The map T is of class $C^{1+\epsilon}$ for some $\epsilon > 0$,
- (3) |T'(x)| > 1 for every $x \in I_1 \cup \cdots \cup I_n$.

The dynamics is concentrated in the repeller Λ of T, which defined by

$$\Lambda := \bigcap_{n=0}^{\infty} T^{-n} I.$$

One of the most important objects in thermodynamic formalism is the so called $topological\ pressure$ of T. It is defined by

(6)
$$P(t) := \sup \left\{ h(\mu) - t \int \log |T'| \ d\mu : \mu \in \mathcal{M} \right\},$$

where \mathcal{M} denotes the space of T-invariant probability measures and $h(\mu)$ the entropy of the measure μ (see [Wa, Chapter 4] for the definition and properties). For each $t \in \mathbb{R}$ there exists a unique measure μ_t attaining the supremum in equation (6), that we call *equilibrium measure* (see [PU, Chapter 5]). It is well known that

the pressure function $t \to P(t)$ is real analytic, strictly convex, strictly decreasing and that $P'(t) = \int \log |T'| d\mu_t$ (see [PP, Chapter 3,4] and [PU, Chapter 5]).

3.1. **Newton map.** A classical result by Bowen [Bo] later generalised by Ruelle [Ru1] establishes a relation between the topological pressure and the Hausdorff dimension of the repeller.

Proposition 3.1 (Bowen's formula). Let T be a cookie cutter map and let $d = \dim_H \Lambda$ then d is the unique root of the equation

$$P(-t\log|T'|) = 0.$$

Let us denote by $N_P(t)$ the Newton map of the pressure function $P(-t \log |T'|)$, that is

$$N_P(t) = t - \frac{P(-t \log |T'|)}{P'(-t \log |T'|)}.$$

We have that $N_P(d) = d$ and the iterations of $N_P(t)$ converge to the unique root of the pressure, that is

$$\lim_{n \to \infty} N_P^n(t) = d.$$

3.2. **Lyapunov spectrum.** The Lyapunov exponent of any given cookie cutter can attain an interval of values $[\alpha_{\min}, \alpha_{\max}]$. For $\alpha \in [\alpha_{\min}, \alpha_{\max}]$ consider,

$$J(\alpha) = \{x \in I : \lambda(x) = \alpha\}$$
 and $J' = \{x \in I : \text{the limit } \lambda(x) \text{ does not exists}\}$.

These level sets induce a decomposition of the repeller Λ of T. Indeed, $\Lambda = J' \cup (\cup_{\alpha} J(\alpha))$. The function that encodes this decomposition is called *Lyapunov* spectrum and it is defined by

$$L(\alpha) := \dim_H(J(\alpha)).$$

H. Weiss [We], building up on previous joint work with Pesin [PW], proved that the Lyapunov spectrum is real analytic. A rather surprising result in light of the fact that this decomposition is fairly complicated. For instance, each level set turns out to be dense in Λ (see [We]). By means of the thermodynamic formalism it is possible to obtain formula for the Lyapunov spectrum. Indeed,

(7)
$$L(\alpha) = \frac{1}{\alpha} \inf_{t \in \mathbb{R}} (P(-t \log |T'|) + t\alpha).$$

This formula follows form the work of Weiss [We] and can be found explicitly, for instance, in the work of Kesseböhmer and Stratmann [KS] (see also [IK] for some explicit computations). Actually, in this setting, the Lyapunov spectrum can be written as

$$L(\alpha) = \frac{1}{\alpha} (P(-t_{\alpha} \log |T'|) + t_{\alpha} \alpha),$$

where t_{α} is the unique real number such that

$$-\frac{d}{dt}P(-t\log|T'|)\Big|_{t=t_{\alpha}} = \alpha.$$

If we denote by μ_{α} the unique equilibrium measure corresponding to the potential $-t_{\alpha} \log |T'|$ we have that

$$\int \log |T'| \ d\mu_{\alpha} = \alpha.$$

Therefore.

(8)
$$L(\alpha) = \frac{1}{\alpha} (P(-t_{\alpha} \log |T'|) + t_{\alpha} \alpha),$$

After the substitution $\alpha = \alpha(t)$, equation (8) becomes:

(9)
$$L(\alpha(t)) = \frac{1}{\alpha(t)} \left(P(-t \log |T'|) + t\alpha(t) \right).$$

4. Main result

In virtue of the comments we have made so far, we can now establish a relation between the Lyapunov spectrum of a cookie cutter map and the Newton map corresponding to the pressure function.

Theorem 4.1. Let T be a cookie cutter map, then for every $t \in \mathbb{R}$ we have

(10)
$$L(-P'(-t\log|T'|)) = N_P(t).$$

Proof. The result follows noticing that the Lyapunov spectrum is given by (see equation (9))

$$L(-P'(-t\log|T'|)) = \frac{1}{-P'(-t\log|T'|)} \left(P(-t\log|T'|) - tP'(-t\log|T'|) \right).$$

This combined with Lemma 5 yields the desired result.

References

- [A] V.I. Arnold, Mathematical methods of classical mechanics. Second edition. Graduate Texts in Mathematics, 60. Springer-Verlag, New York, 1989. xvi+508 pp.
- [Ba] L. Barreira Dimension and recurrence in hyperbolic dynamics. Progress in Mathematics, 272. Birkhuser Verlag, Basel, 2008. xiv+300 pp.
- [Bo] R. Bowen, Hausdorff dimension of quasi-circles. Inst. Hautes Etudes Sci. Publ. Math. 50, 11-26 (1979).
- [Fa] K. Falconer Fractal geometry. Mathematical foundations and applications. Second edition. John Wiley and Sons, Inc., Hoboken, NJ, 2003.
- [IK] G. Iommi and J. Kiwi The Lyapunov spectrum is not always concave Journal of Statistical Physics vol.135 no.3, 535-546 (2009).
- [K] G. Keller, Equilibrium states in ergodic theory, London Mathematical Society Student Texts, 42. Cambridge University Press, Cambridge, 1998.
- [KS] M. Kesseböhmer and B. Stratmann A multifractal analysis for Stern-Brocot intervals, continued fractions and Diophantine growth rates Journal für die reine und angewandte Mathematik (Crelles Journal) 605, 133-163, (2007).
- [Mi] J. Milnor Dynamics in one complex variable. Third edition. Annals of Mathematics Studies, 160. Princeton University Press, Princeton, NJ, 2006. viii+304 pp.
- [PP] W. Parry and M. Pollicott, Zeta functions and the periodic orbit structure of hyperbolic dynamics. Astrisque No. 187-188 (1990), 268 pp.
- [PW] Y. Pesin, and H. Weiss A multifractal analysis of equilibrium measures for conformal expanding maps and Moran-like geometric constructions. J. Statist. Phys. 86 (1997), no. 1-2, 233–275
- [Pe] Y. Pesin Dimension Theory in Dynamical Systems CUP (1997).
- [PU] F. Przytycki and M. Urbański, Conformal Fractals: Ergodic Theory Methods, Cambridge University Press 2010.
- $[\mathrm{Ru1}]$ D. Ruelle Repellers for real analytic maps. Ergodic Theory and Dynamical Systems 2 (1982), no. 1, 99107.
- [Ru2] D. Ruelle, Thermodynamic formalism. The mathematical structures of equilibrium statistical mechanics, Second edition. Cambridge Mathematical Library. Cambridge University Press, Cambridge, 2004.

- [Wa] P. Walters, An introduction to ergodic theory, Graduate Texts in Mathematics 79. Springer-Verlag, New York-Berlin, 1982.
- [We] H. Weiss The Lyapunov spectrum for conformal expanding maps and axiom-A surface diffeomorphisms. J. Statist. Phys. 95 no. 3-4, , 615–632 (1999).

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